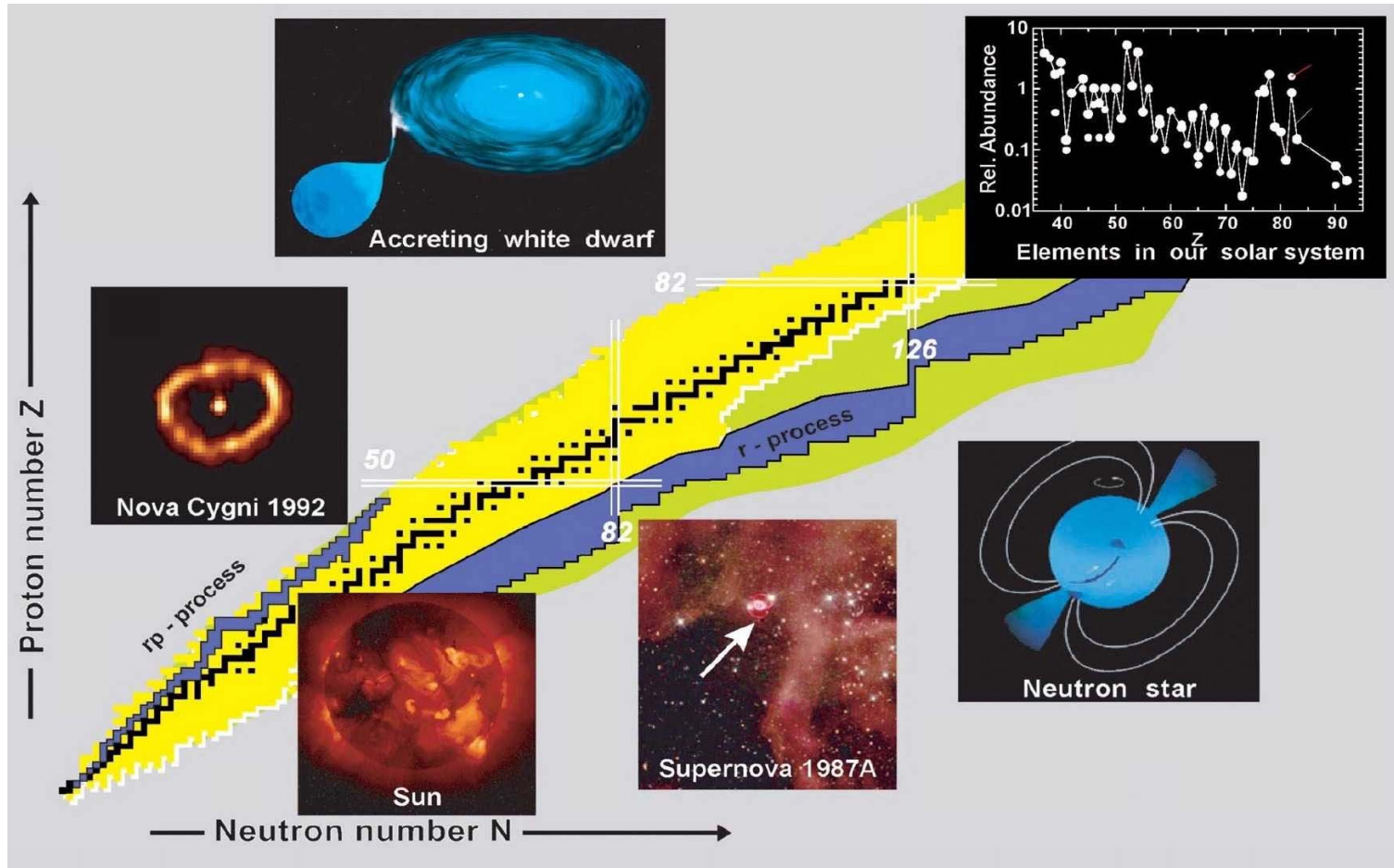


K.-L. Kratz
 Kernchemie + VISTARS, Universität Mainz

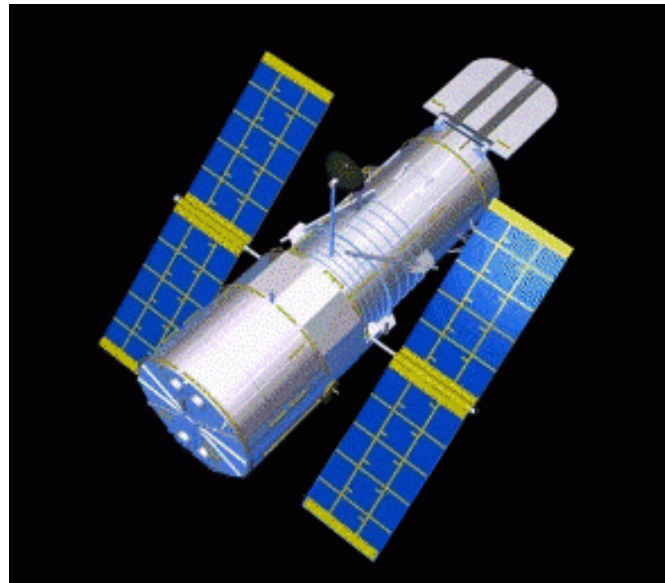


Historically, nuclear astrophysics has been concerned with

- Interpretation of observed **abundance distributions**, or specific signatures of **luminosity curves**;
- Description of originating **nucleosynthesis processes**
(see, e.g. B²FH, 1957!)

Observational instrumentation:

- **meteoric** and overall **solar** abundances;
- ground- and satellite-based telescopes, like *Imaging Spectrograph* (STIS) at **Hubble** or *HIRES* at **Keck I**, and γ -ray satellites like *INTEGRAL* or x-ray observatories *CHANDRA* and *XMM-Newton*



Nuclear Physics in Stellar Binaries

Cataclysmic binaries are potential site for explosive nuclear reaction driven processes
- thermonuclear runaway -



- Nova: thermonuclear runaway burning on accreting white dwarfs
- Supernova type I: cataclysmic burning on white dwarfs
- X-ray bursts: thermonuclear runaway on accreting neutron stars
- X-ray pulsars: steady burning on accreting neutron stars

rp-process

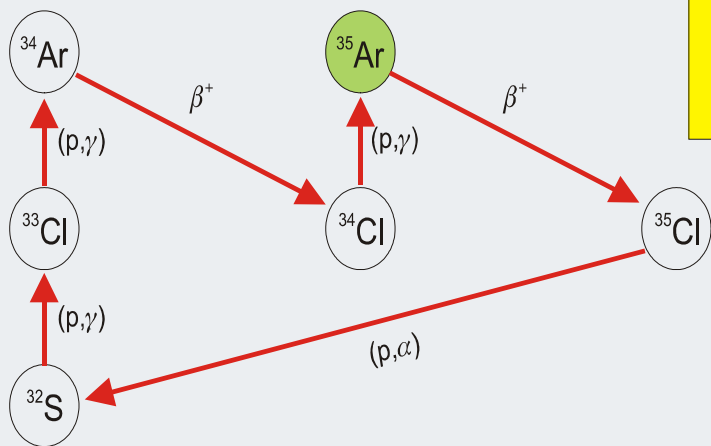
at "low" temperatures

Novae

explosive H-burning on surface of white dwarfs

cycle pattern due to Cb-barriers;
p-capture and β^+ -decay
with similar time scales;
path close to stability

up to ${}_{19}\text{K} - {}_{23}\text{V}$
(p, γ); (γ ,p); (p, α); $T_{1/2}(\beta^+)$
on odd-Z, $T_z = +\frac{1}{2}$ nuclei
e.g. ${}^{23}\text{Mg}$, ${}^{27}\text{Si}$, ${}^{31}\text{S}$, ${}^{43}\text{Ti}$



nucl. masses
reaction Q-values

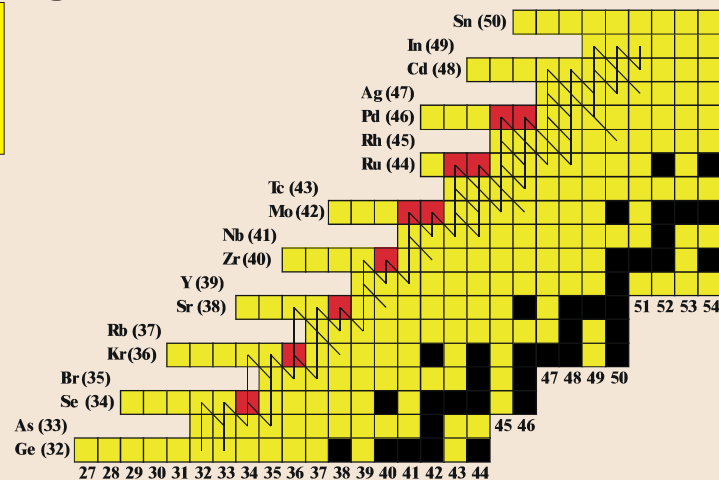
at "high" temperatures

X-ray bursts

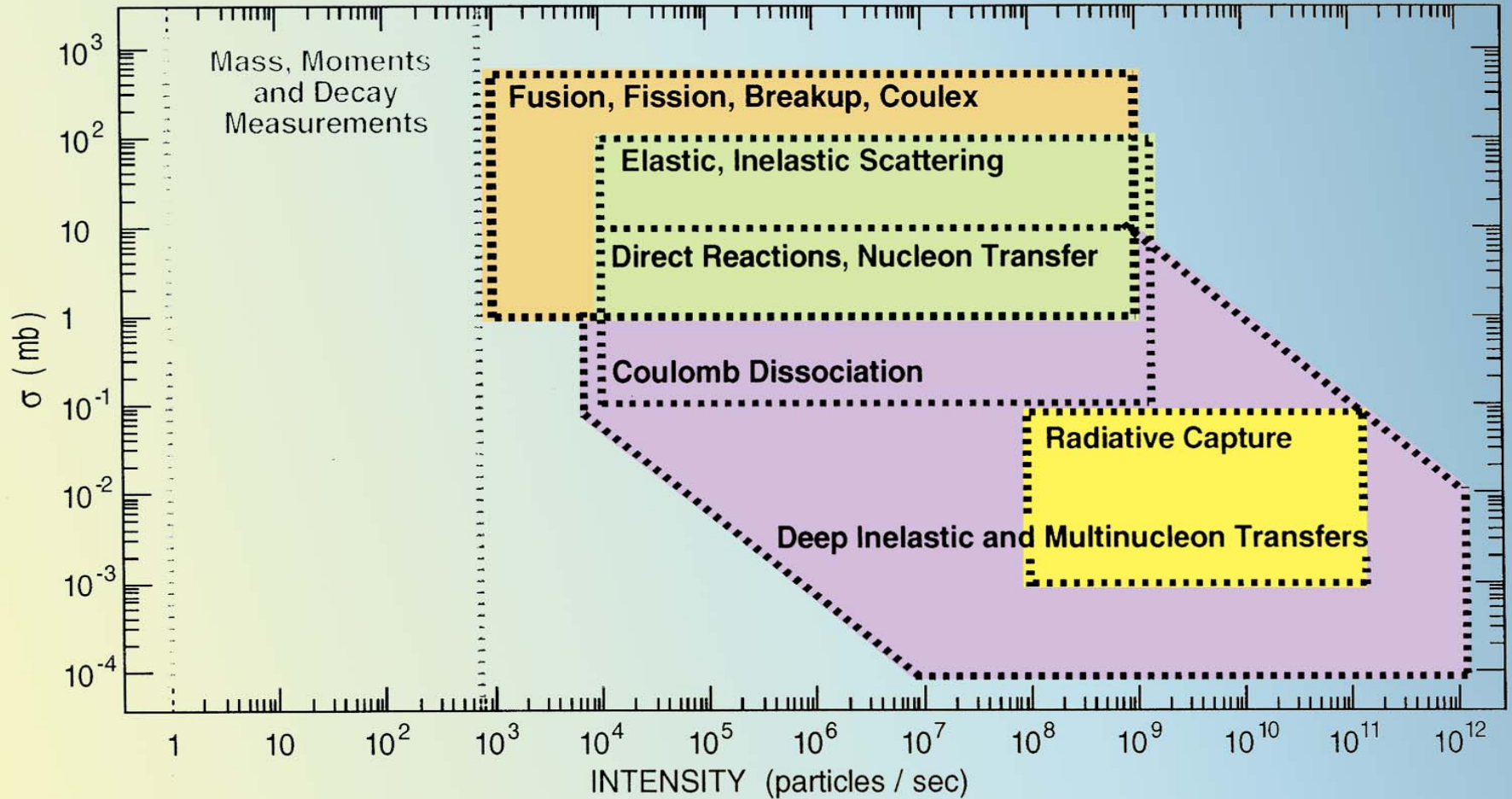
explosive H- and He-burning on surface of accreting
neutron stars

(p, γ)-(γ ,p) equilibrium "waiting-point" concept;
sequence of fast p-captures
and subsequent β^+ -decay
path far away from stability

up to ${}_{46}\text{Pd} - {}_{50}\text{Sn}$
(p, γ); (γ ,p); (2p, γ); $T_{1/2}(\beta^+)$
on even Z, $T_z = -\frac{1}{2}$ nuclei
e.g. ${}^{68}\text{Se}$, ${}^{72}\text{Kr}$, ${}^{76}\text{Sr}$, ${}^{80}\text{Zr}$, ${}^{84}\text{Mo}$...

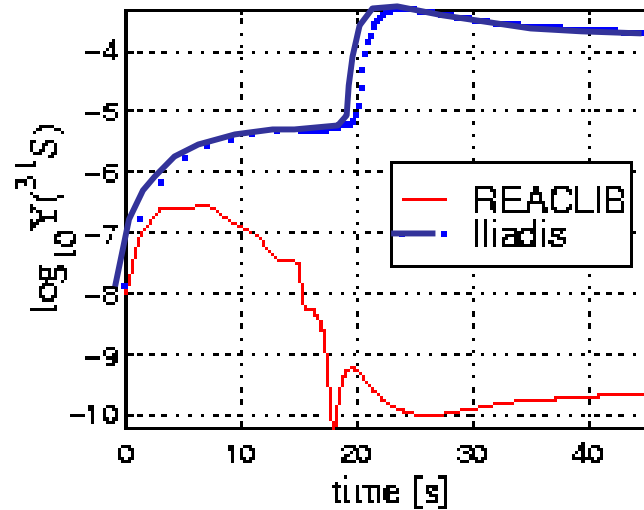
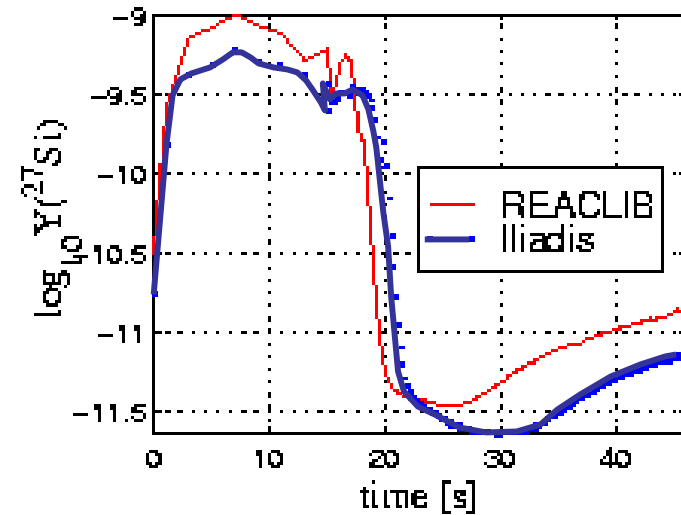


Most relevant Nuclear Structure Experiments

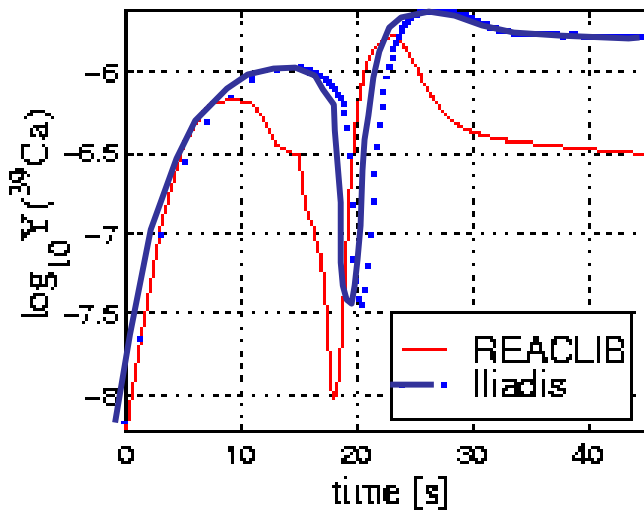
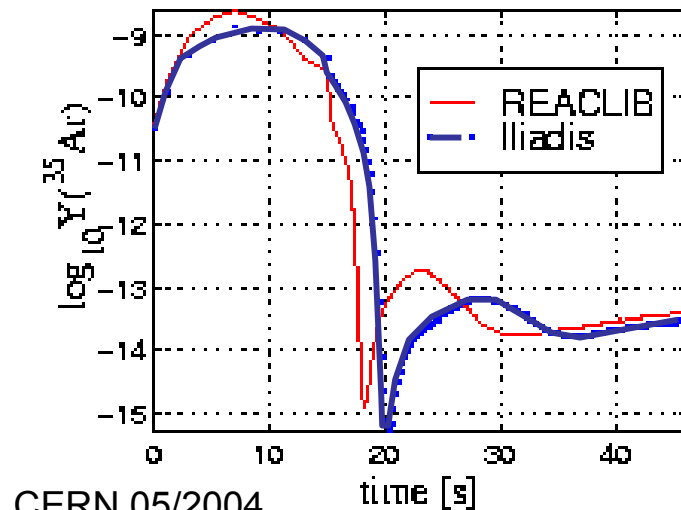


Self consistent calculations

Consequences for isotopic abundances!



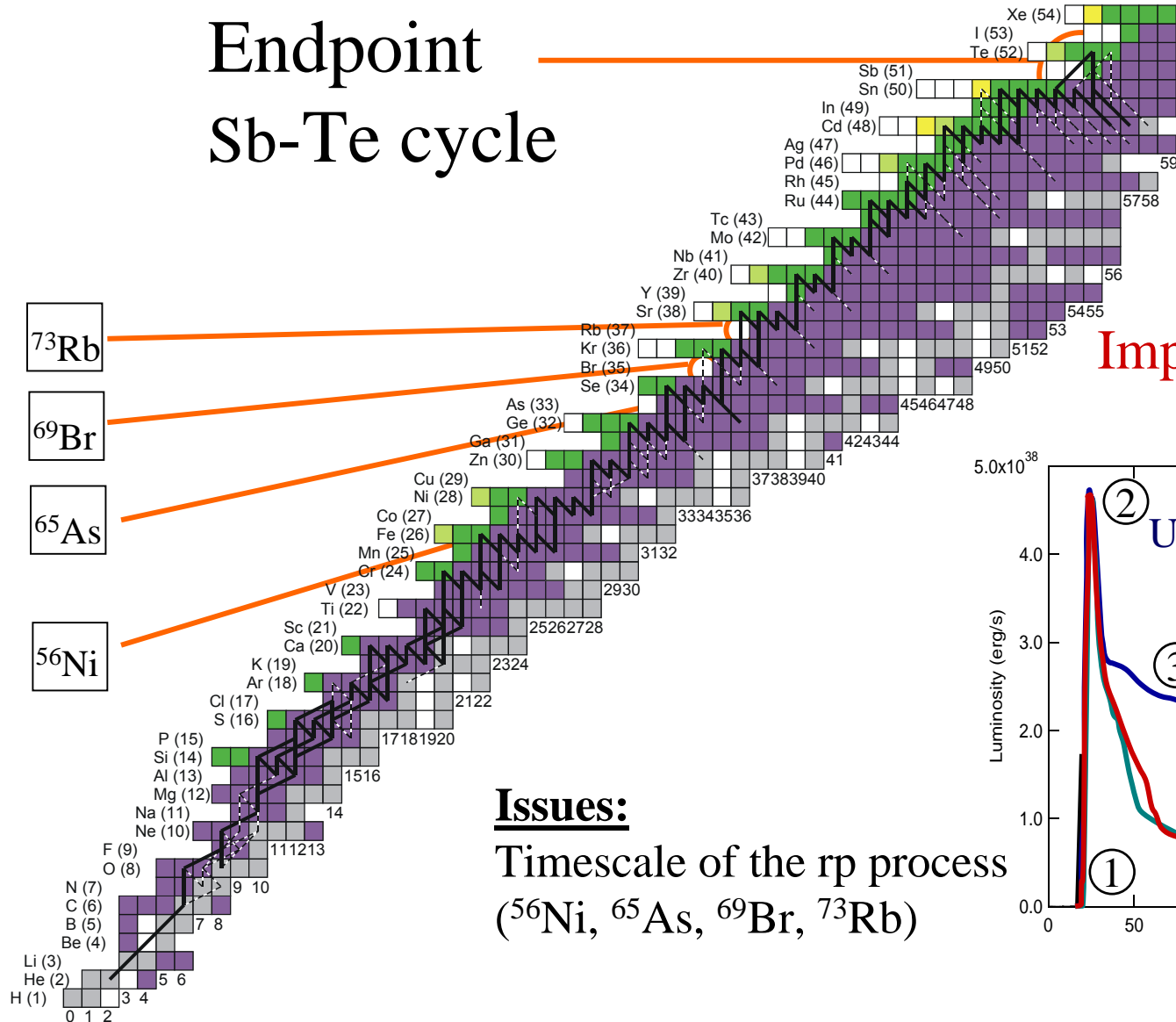
^{31}S , slower depletion
 \Rightarrow mixed towards surface



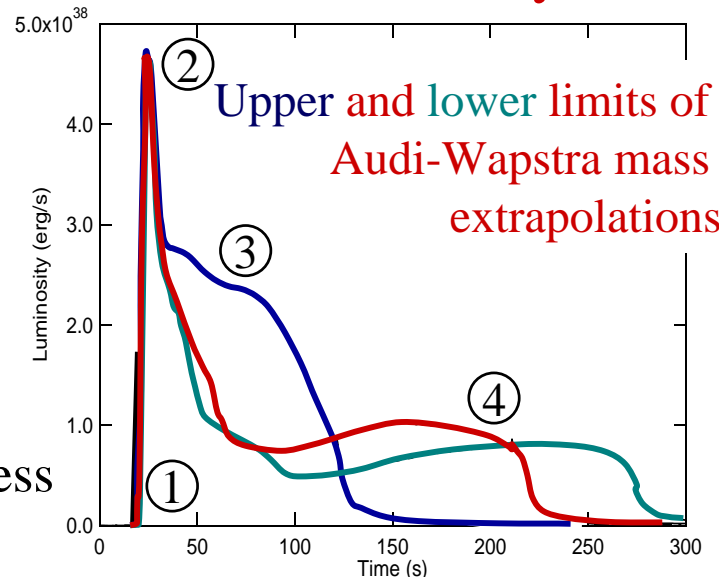
^{39}Ca , slower depletion
 \Rightarrow mixed towards surface

Waiting points in the rp-process

Endpoint
Sb-Te cycle



Impact on X-ray burst
luminosity curve



Issues:

Timescale of the rp process
(^{56}Ni , ^{65}As , ^{69}Br , ^{73}Rb)

from Wiescher et al. (2003)

The Energy production in X-ray bursts

characterized by 4 periods at increasing temperature ($2 \times 10^8 - 2 \times 10^9$ K):

① ignition phase

Hot CNO-cycles, triggered by p-captures on C, N, O isotopes; formation of ^{14}N , $^{14,15}\text{O}$, ^{17}F and ^{18}Ne waiting-point nuclei.

② ignition of triple- α -process

Rapid depletion of $^{14,15}\text{O}$ and ^{18}Ne by α -reactions to ^{24}Si ; further p- and α -captures to waiting-point isotopes ^{29}S , ^{34}Ar ; and subsequent α -captures to **main waiting-point nucleus ^{56}Ni** .

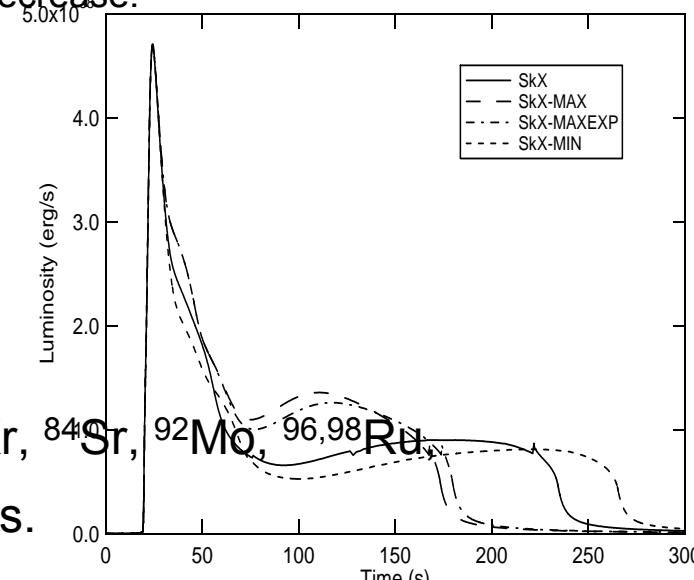
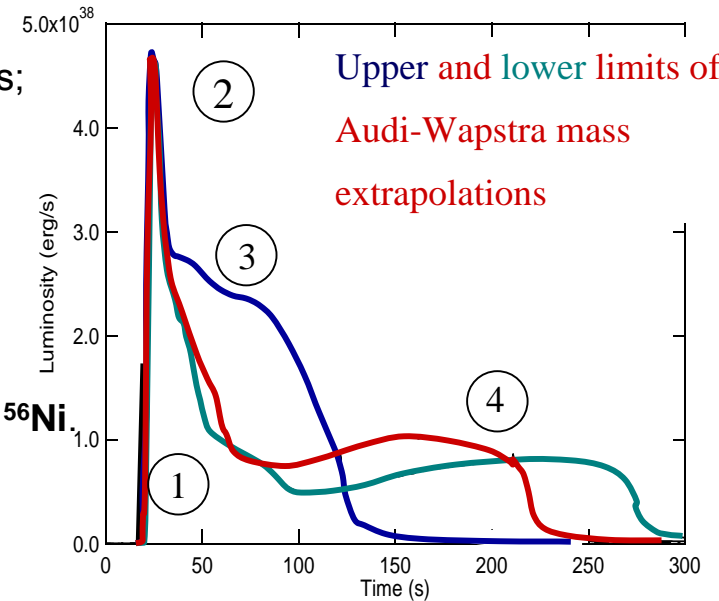
③ dormant period

Here, further processing is halted by (p, γ)-(γ ,p) equilibrium; energy production drops rapidly; and temperature starts to decrease.

④ rp-process beyond ^{56}Ni

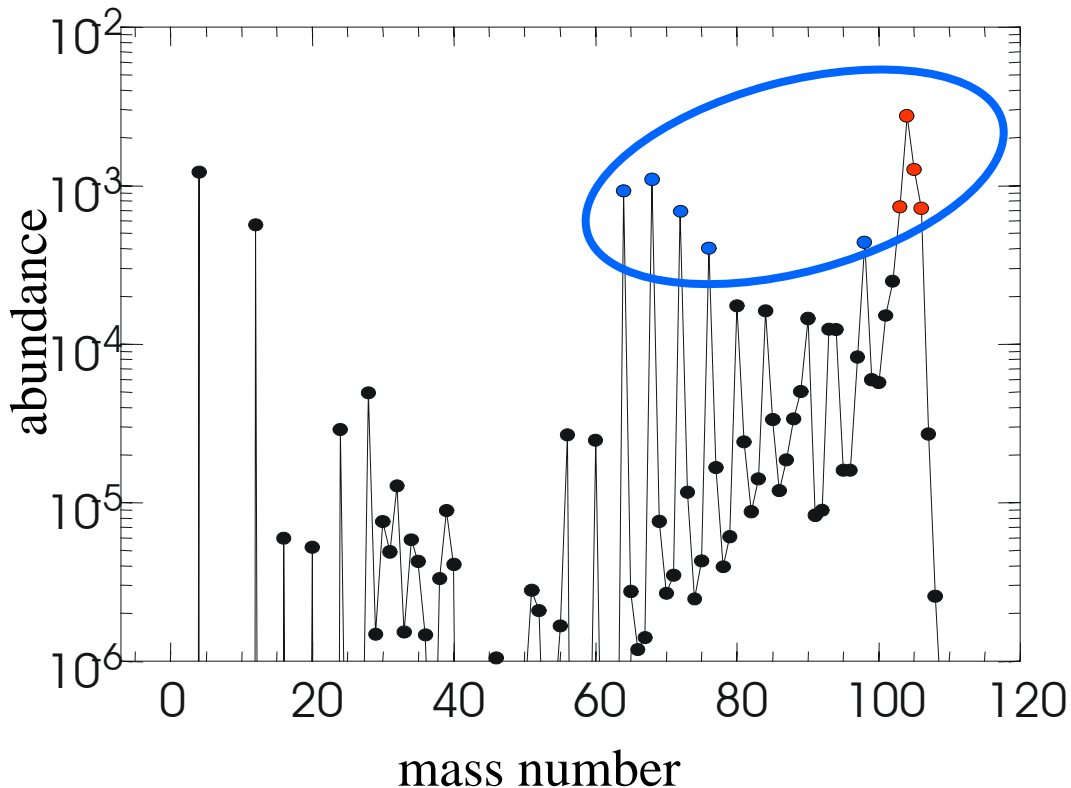
Start at $T_g = 1.5 \times 10^9 - 1.0 \times 10^9$ K; cooling phase down to $\sim 5 \times 10^8$ K; several even-even waiting-points ^{64}Ge , ^{68}Se , ... ^{100}Sn ; overall time scale ~ 500 s.

Clear abundance peaking of light **p-nuclei** ^{74}Se , ^{78}Kr , ^{84}Sr , ^{92}Mo , $^{96,98}\text{Ru}$, so far not explained in classical p-process scenarios.



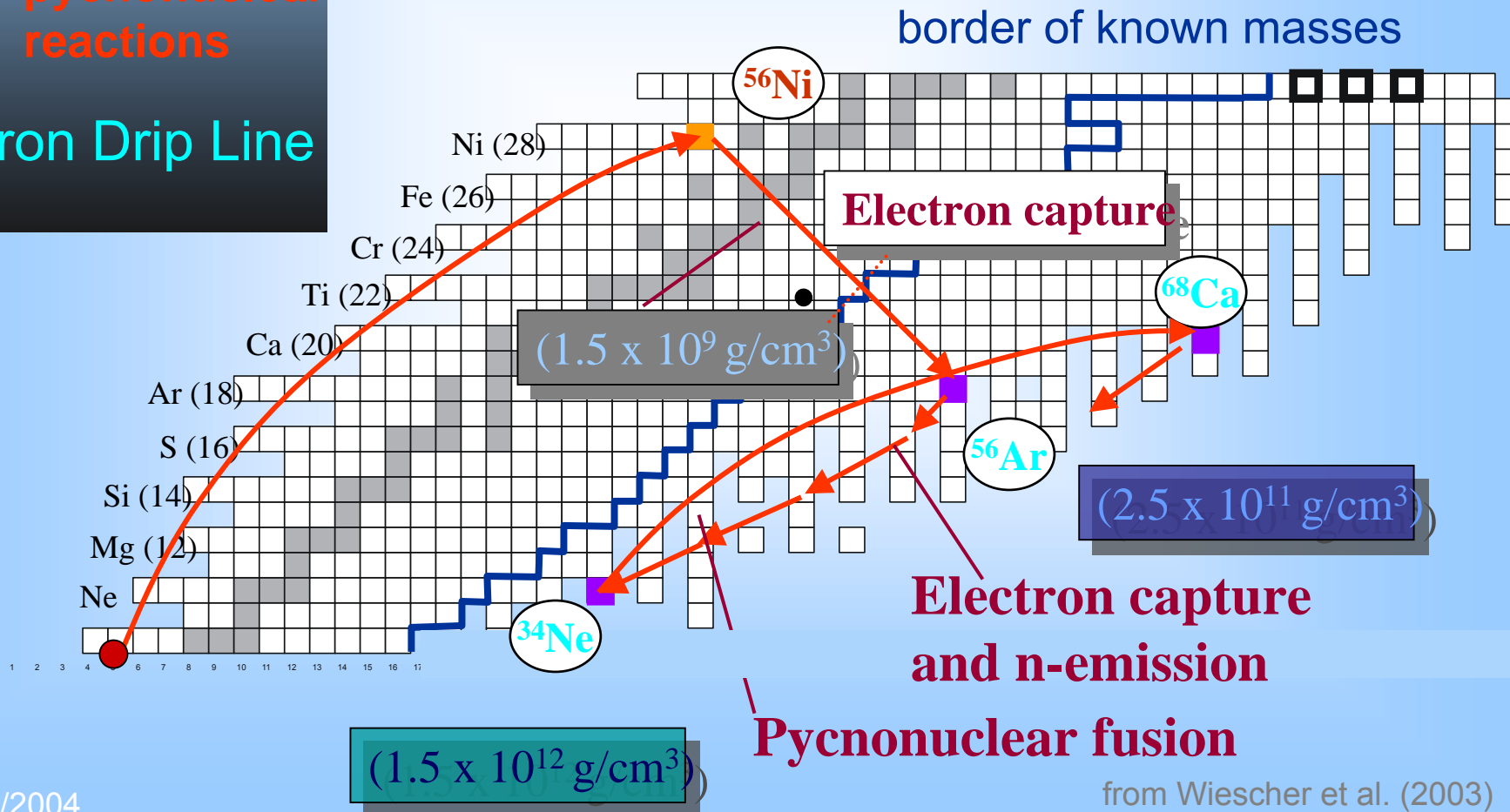
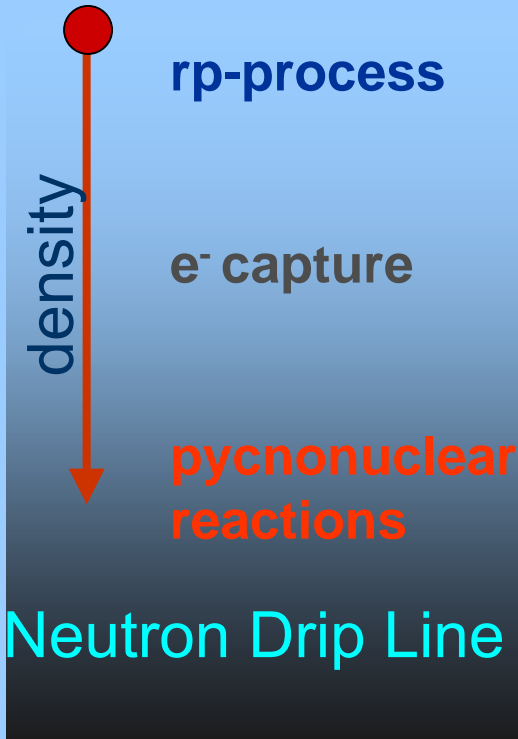
The fate of the processed matter?

Final abundance distribution



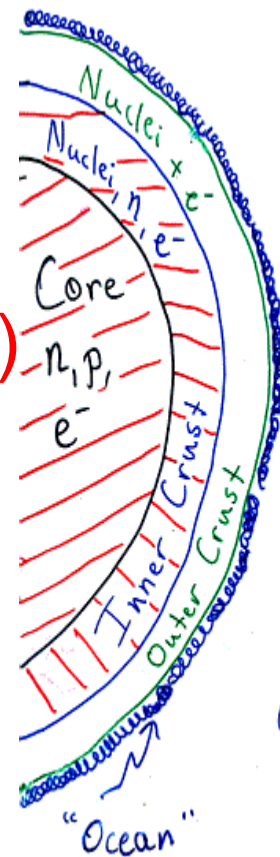
- Ejected out of neutron star grav. potential
⇒ contribution to light p-nuclei abundances
- Embedded into neutron star crust
⇒ modification of crust composition

The fate of matter in the neutron star crust



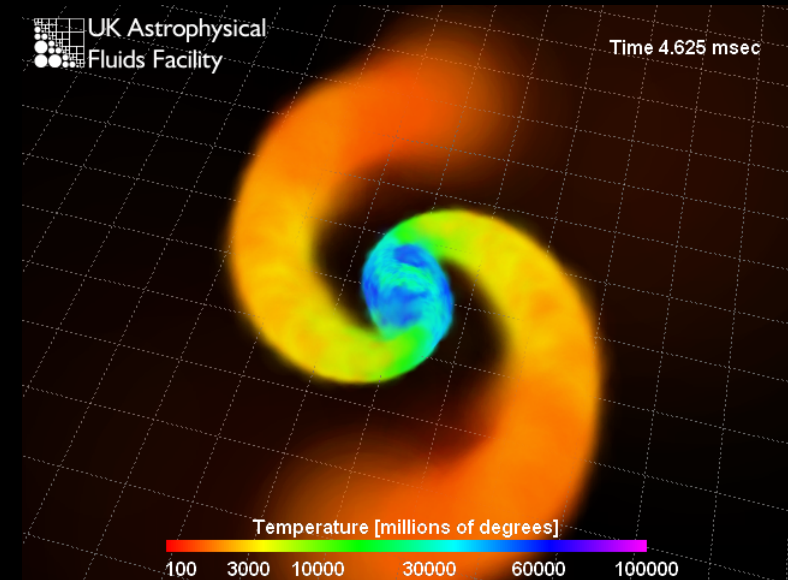
Conclusion

- Nuclear structure along rp-process reaction path ($N=Z$ line) determines abundance distribution in the ashes
- Endpoint is associated with α -unbound nuclei beyond $Z, N=50$
- **Electron capture drives abundance distribution towards the neutron drip line (timescale $>10^5$ y)**
- Electron capture induced neutron emission drives abundance distribution towards $Z < 12$
- Pycno-nuclear reactions



Nuclear Physics in Neutron-Rich Explosive Environments

High neutron to seed ratio (10 - 150) for rapid neutron-capture processes only in explosive astrophysical sites



- Supernova type II: explosive shell burning (He, C)
high-entropy neutrino wind ejecta
- Neutron-star (NS) mergers: low-entropy ejecta of NS matter
- Axial jets in SNe: blast ejecta from rotating NS accretion disk

r-process

Supernova II

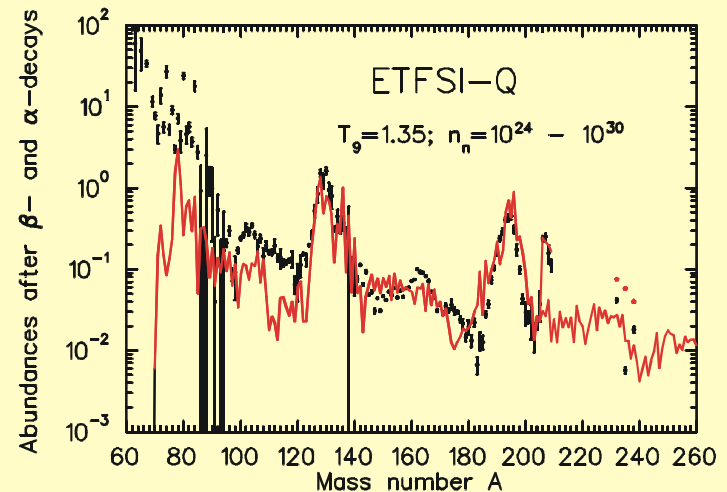
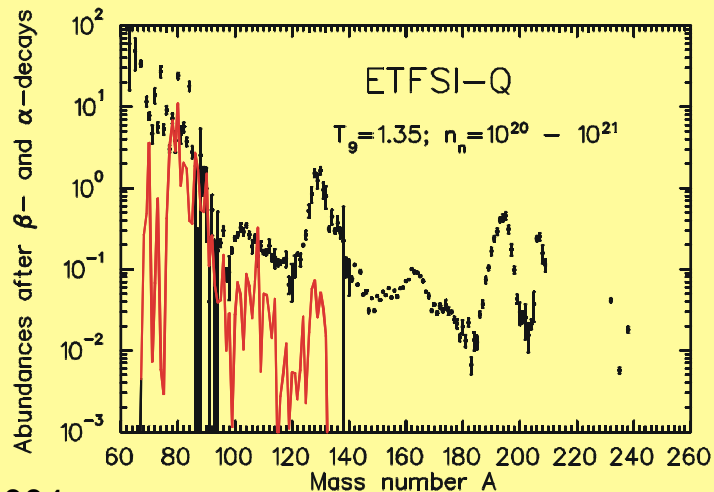
moderate n-densities ($n_n \leq 10^{23}$)
explosive (He, C, Si) shell burning
high-entropy bubble, core scenario

Neutron Star Mergers

high n-densities ($n_n \approx 10^{30}$)
very n-rich ejecta
fission “recycling”

nuclear data needs

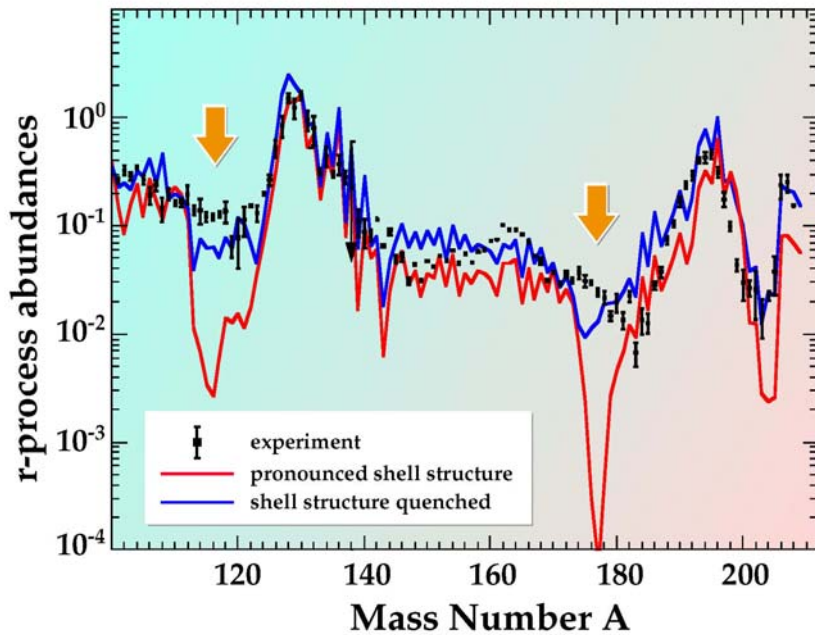
direct: Q_β , S_n , $T_{1/2}$, P_n , $\sigma_{n\gamma}$ of “waiting-point nuclei” at N_{mag}
indirect: development of nuclear structure with isospin



New Nuclear structure at drip-lines, as „shell quenching“?

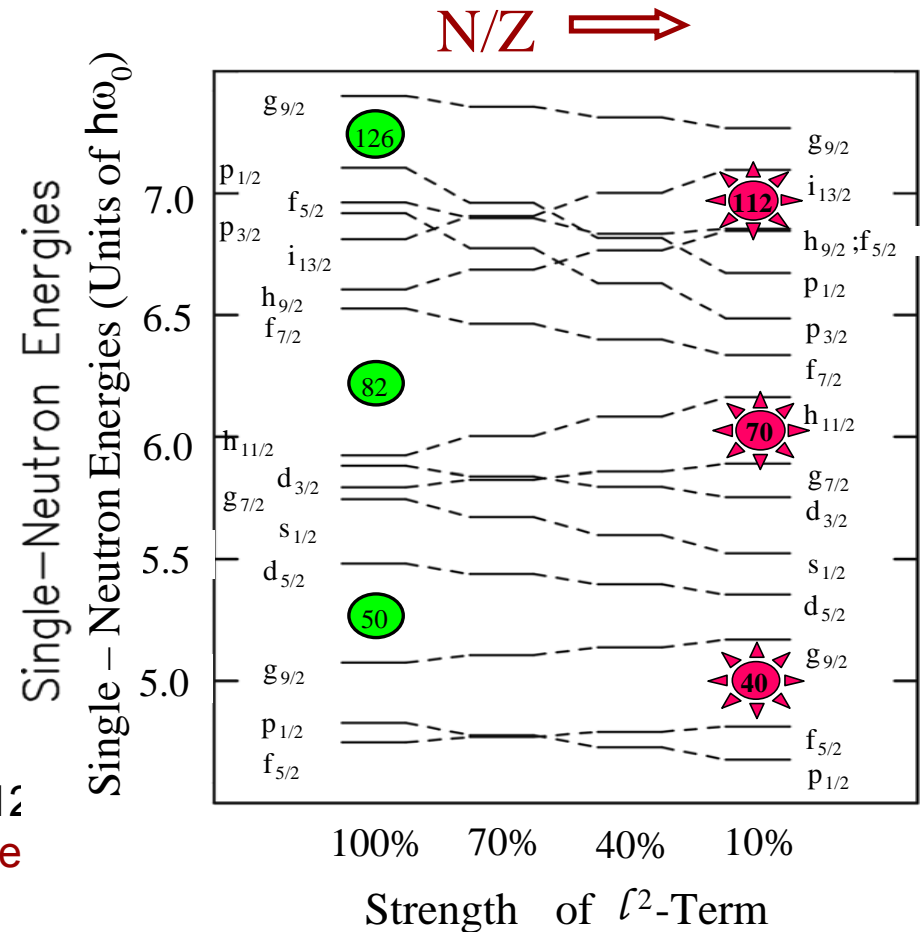
Deficiencies prior to the main peaks were attributed by our group to nuclear structure effects:

- too strong shell strength for extremely neutron-rich magic nuclei far from stability
- nuclear models adjusted to stable nuclides
- new physics far from stability?



FK²L (Ap.J. 403 ; 1993)

“..the calculated r-abundance ‘hole’ in the $A \approx 112$ region reflects ... **the weakening of the shell structure** ... below ^{132}Sn “
50 82

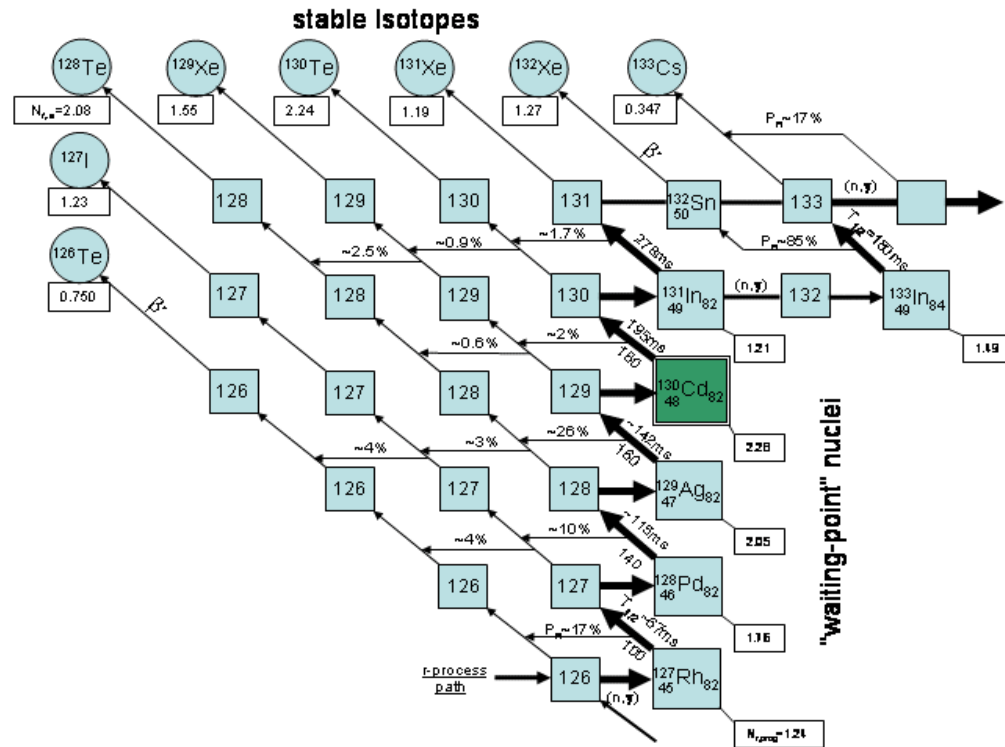
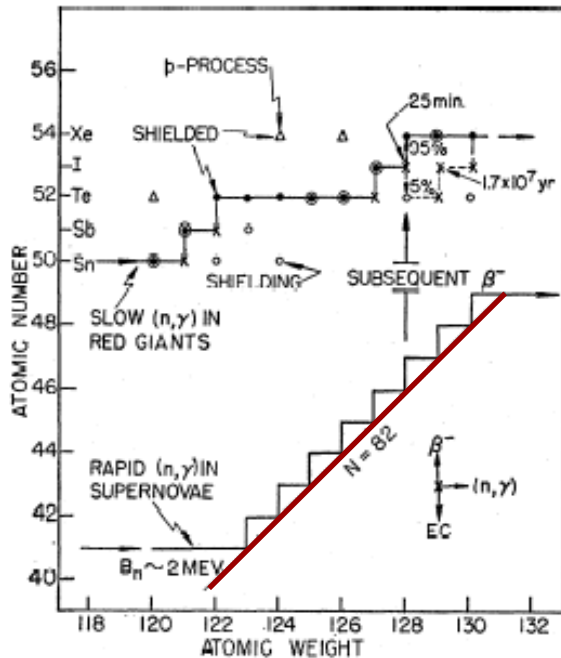


B. Pfeiffer et al.,
Acta Phys. Polon. **B27** (1996)

R-abundance peaks and neutron-shell numbers

already B²FH (Revs. Mod. Phys. 29; 1957)
 C.D. Coryell (J. Chem. Educ. 38; 1961)

...still today important r-process properties to be studied experimentally and theoretically.



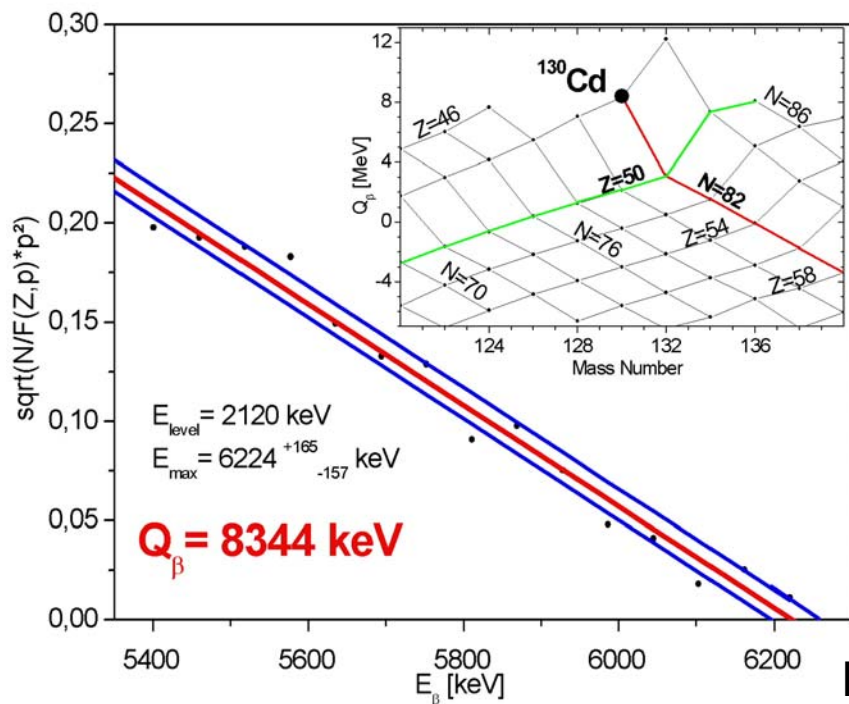
“climb up the staircase” at N=82;
 major waiting point nuclei;
 “break-through pair” ¹³¹In, ¹³³In;

K.-L. Kratz (Revs. Mod. Astr. 1; 1988)
 climb up the N= 82 ladder ...
 A ≅ 130 “bottle neck”

↻ “association with the rising side of major peaks in the abundance curve”

⇒ **total r-process duration τ_r**

Q_β -value of $^{130}\text{Cd}_{82}$



Way-Wood diagram

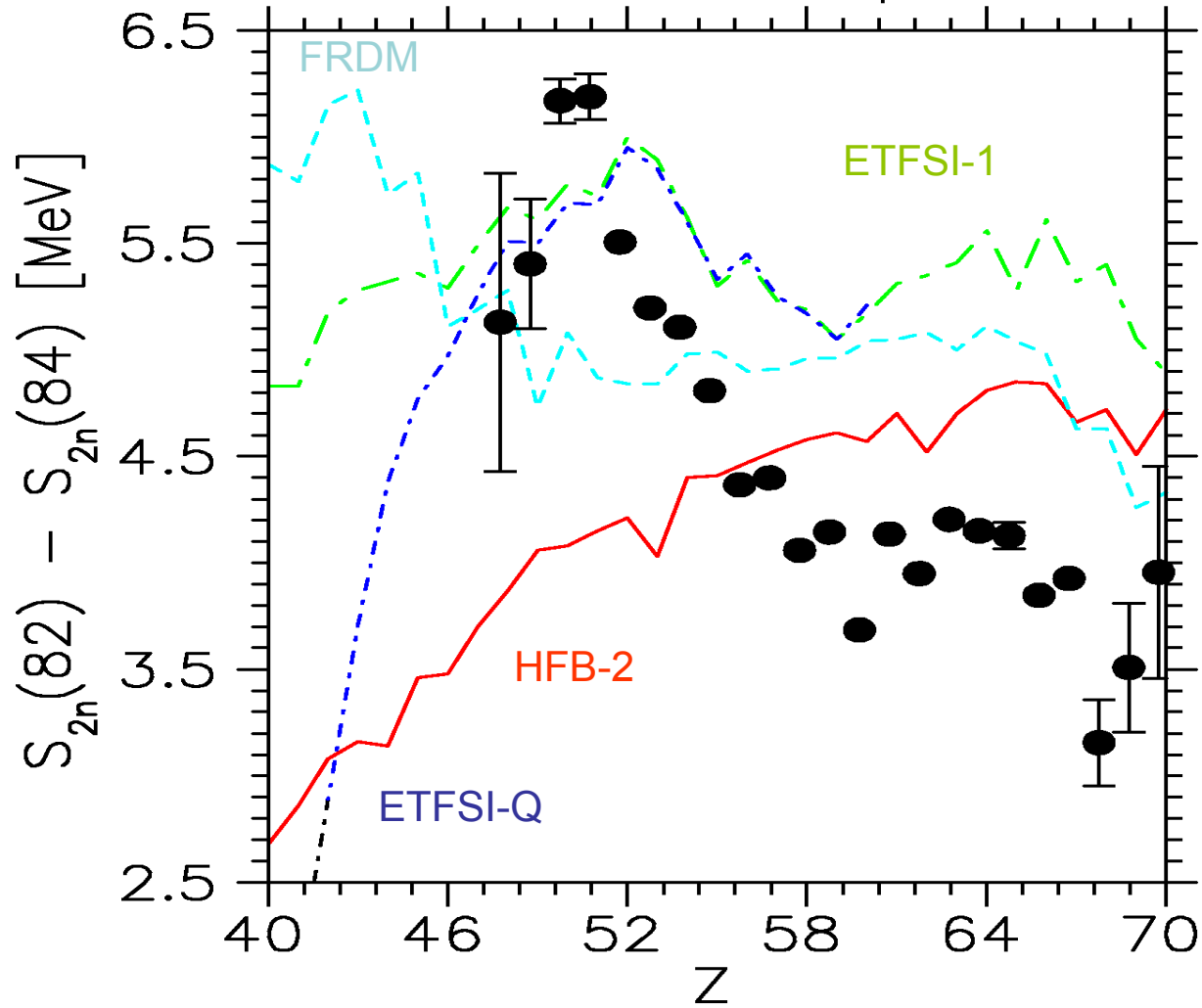
⇒ **Z=50** and **N=82** shell closures „visible“

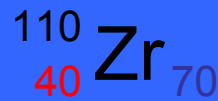
High Q_β -value is a clear signature for an $N=82$ „shell quenching“ below $^{132}\text{Sn}_{50}$

Mass model predictions

Mass model predictions	Q_β
Hilf et al. (<i>GTNM</i> , 1976)	7.57 MeV
Möller et al. (<i>FRDM</i> , 1995):	7.43 MeV
Aboussir et al. (<i>ETFSI</i> , 1995):	7.87 MeV
Duflo & Zuker (1995)	7.56 MeV
Dobaczewski et al. (<i>HFB/SkP</i> , 1996):	8.93 MeV
Pearson et al. (<i>ETFSI-Q</i> , 1996):	8.30 MeV
Audi & Wapstra (<i>Mass Eval.</i> , 1997):	8.50 MeV
Goriely et al. (<i>HFB/CS</i> , 2001)	7.00 MeV
Samyn et al. (<i>HFB-2</i> , 2002)	7.64 MeV
Brown et al. (<i>local OXBASH</i> , 2003):	8.75 MeV

N=82 Shell Gap



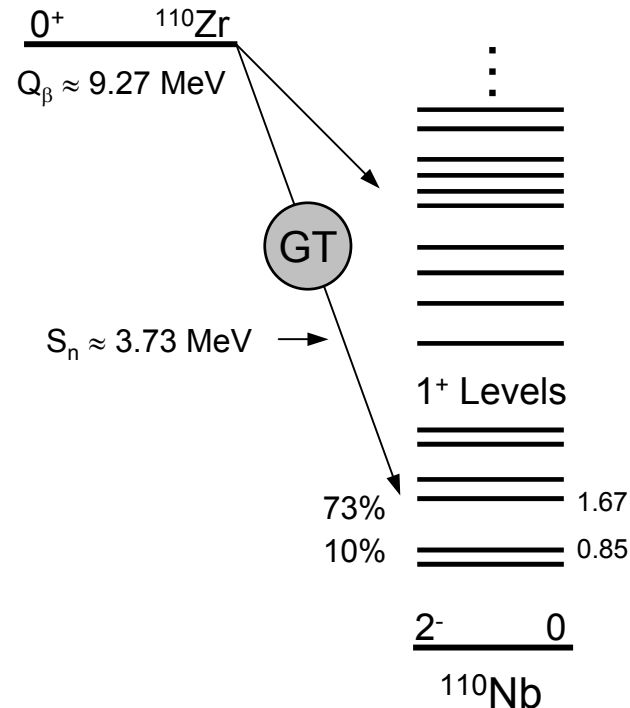


... a new double-magic „waiting-point“ nucleus?

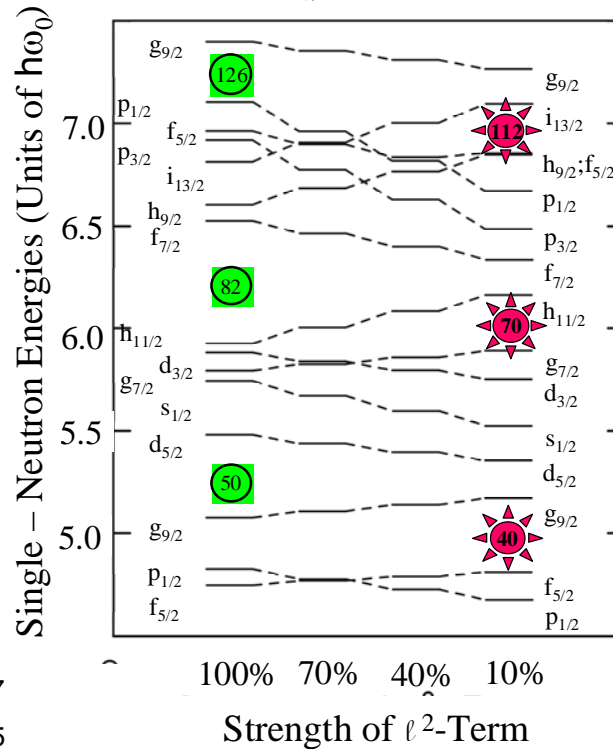
Beta-decay

Normal shell strength
strongly deformed ($\epsilon_2=0.31'$)

$T_{1/2} = 88 \text{ ms}$
 $P_n = 8 \%$

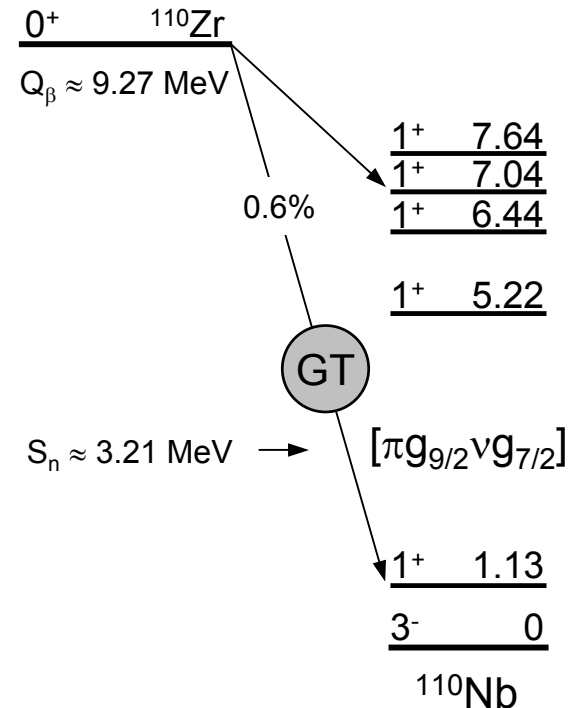


$N/Z \Rightarrow$

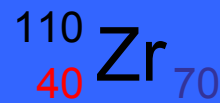


Shell strength quenched
spherical

$T_{1/2} = 14 \text{ ms}$
 $P_n = 0.7 \%$

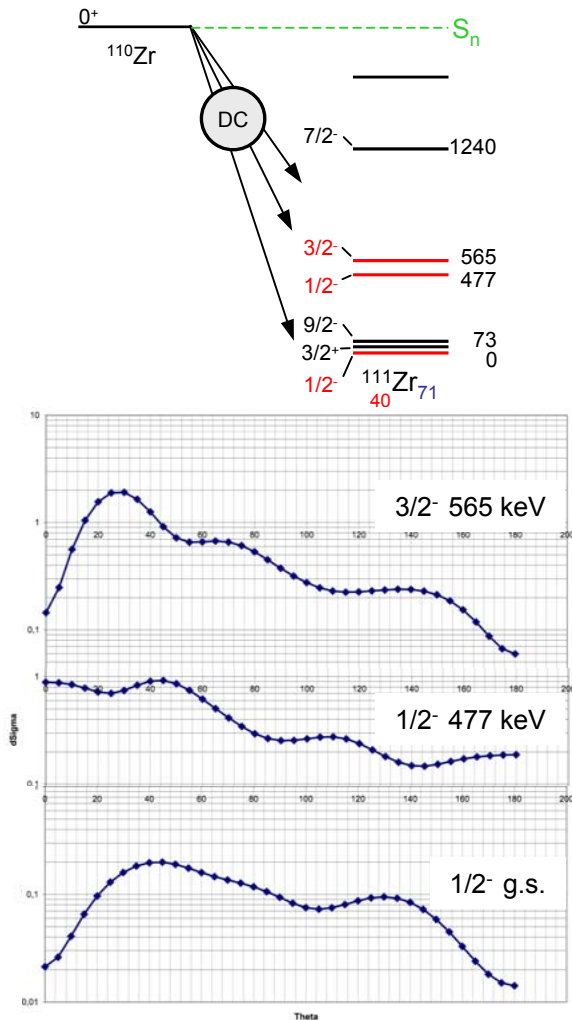


K.-L. Kratz, A. Ostrowski (2004)

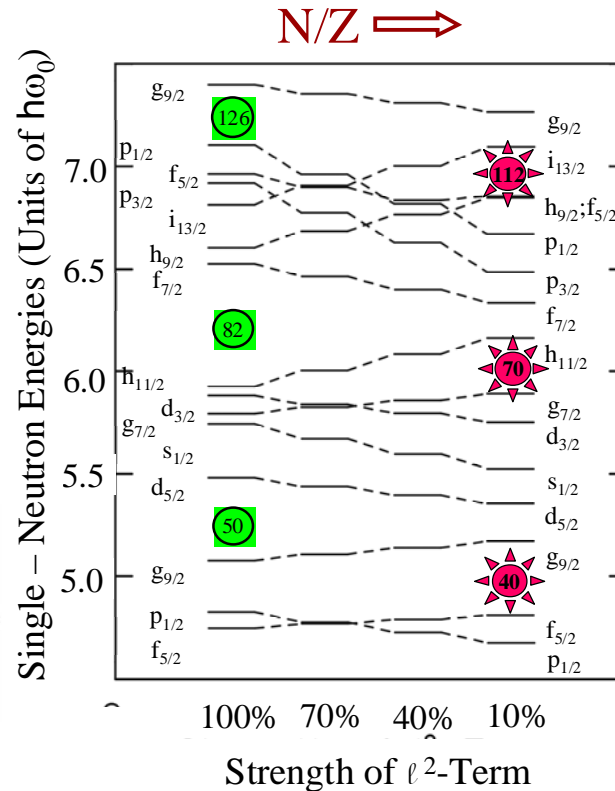


... a new double-magic „waiting-point“ nucleus?

Normal shell strength
strongly deformed ($\epsilon_2=0.34$)



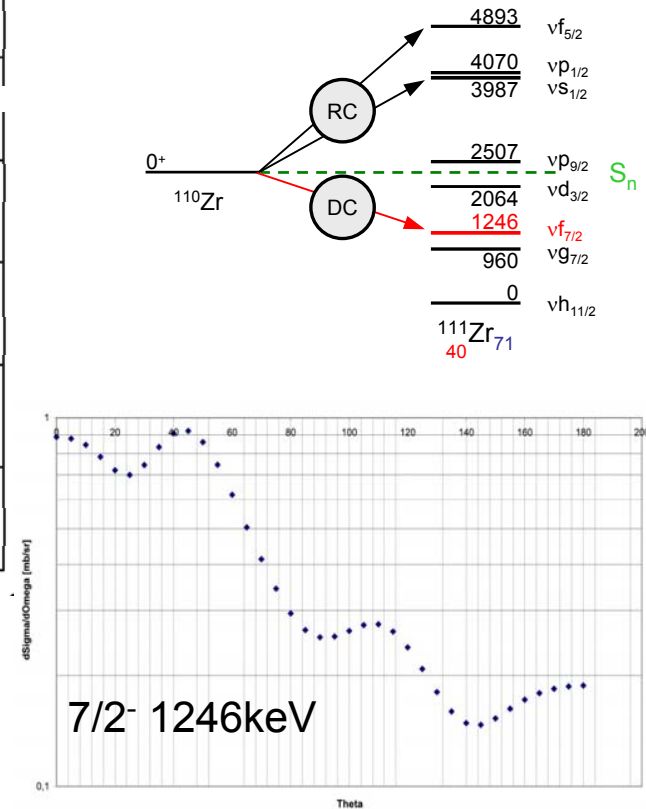
Neutron-capture
via (d,p) in inverse kinematics



B. Pfeiffer et al.,
Acta Phys. Polon. **B27** (1996)

Shell strength quenched
spherical

$$S_n(^{111}\text{Zr}) \approx 2.30 \text{ MeV}$$

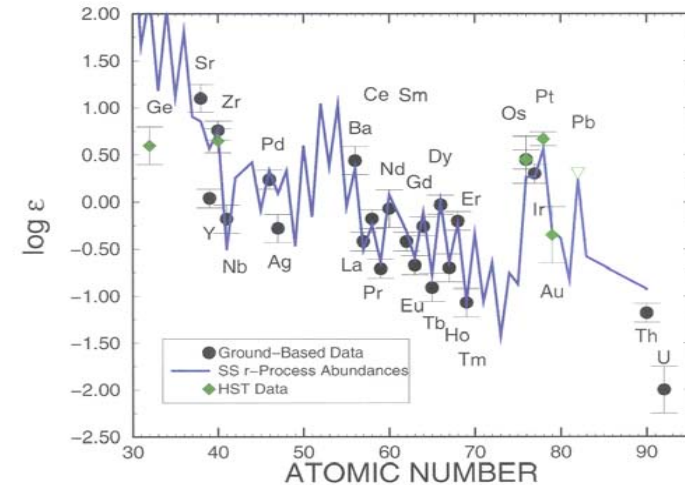


K.-L. Kratz, A. Ostrowski (2004)

Ultra-metal-poor Halo stars

Th-U cosmochronometer

- Th-U chronometer ideal for dating of solar system.
- Lower limit for age of the Universe requires models of Galactic chemical evolution.
- High-resolution optical spectroscopy of ultra-metal-poor, very old Halo red giant stars opens new perspectives:
 - One (or few) nucleosynthesis events seeded ISM.
 - Scaled solar system r-process abundances for $56 \leq Z \leq 79$.
 - Radioactive dating requires „production ratios“.
 - Th/U ratio (hopefully) less affected by extrapolations of nuclear structure

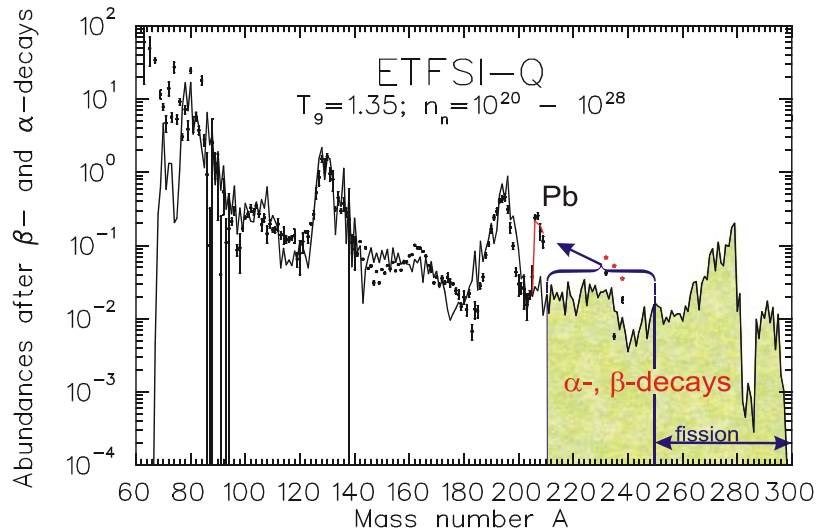


Abundances in BD+17°3248

Calculated abundances of the long-lived actinides test for short-lived isotopes expected in Galactic cosmic rays.

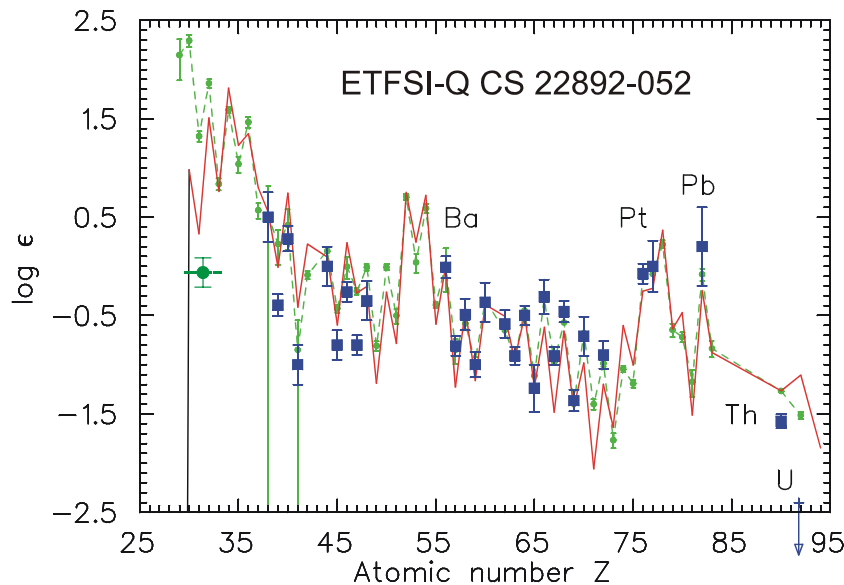
What do we calculate?

Reproduce **isotopic** $N_{r,\odot}$ distribution
using waiting-point approach
(lsq-fits; $A \geq 80$, $A \geq 120$)



Pb abundances as
reliability criterion
for Th, U region!

Convert to elemental $N_{r,\odot}$ curve(—), and scale to halo abundances (■)



↪ comparison to
 $N_{r,\odot}(Z)$

⇒ excellent
agreement above
Ba

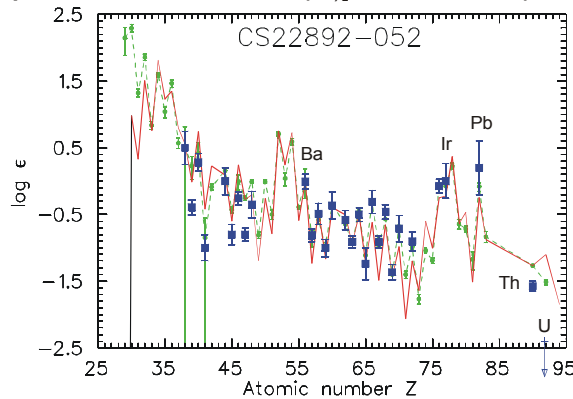
⇒ Th, U cosmo-
chronometry

Cosmochronology with Thorium

Several UMP halo stars consistently show solar elemental abundance pattern, among them

CS 22892-052

with 17 elements between Ba and Ir, plus Pb and ^{232}Th ($T_{1/2} = 14 \cdot 10^9 \text{ a}$)



⇒ age from average of Th/Ba-Ir

$14.6 \pm 3.0 \text{ Gyr}$

≈ age of the Universe



Tuning the clock with Uranium

Because of shorter half-life ($T_{1/2} = 4.5 \cdot 10^9 \text{ a}$), ^{238}U may yield more precise age.

Recently, first observation of U in **CS 31082-001**

(Cayrel *et al.*, 2001)

⇒ age from U/Th

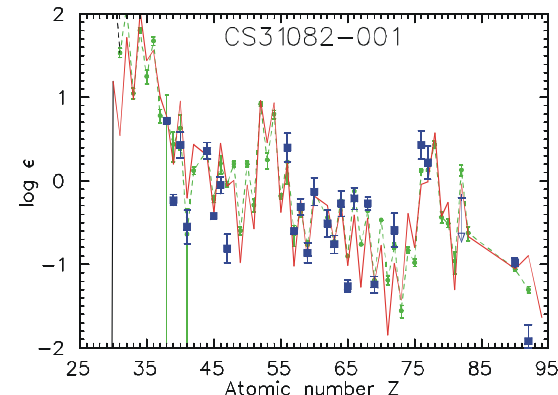
$15.0 \pm 3.0 \text{ Gyr}$

(Schatz *et al.*, 2002)

In the meantime, second observation of U in **BD +17°3248** ⇒ age from U/Th and U,Th/3rd peak

$13.8 \pm 4.0 \text{ Gyr}$

(Cowan *et al.*, 2002)



Isotope ratios in metal-poor, r-process-rich stars

For a limited number of elements (as Rb, Ba, Eu), not only **elemental** but also **isotopic** abundances have been observed in stars.

This provides crucial information on the timing of the neutron-capture processes, the s- and r-process, in the early Galaxy.

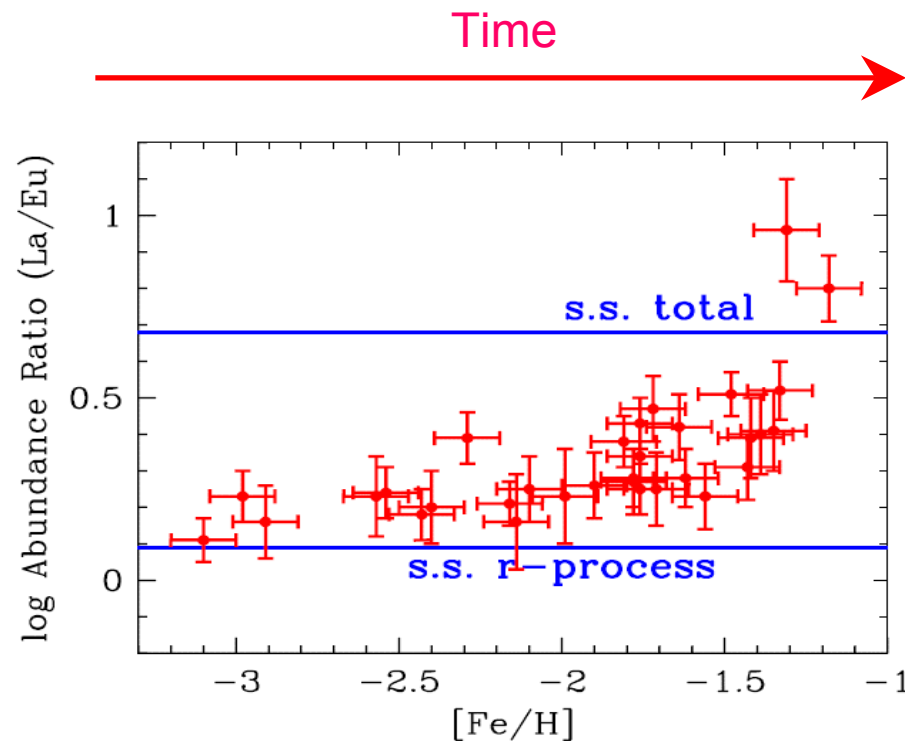
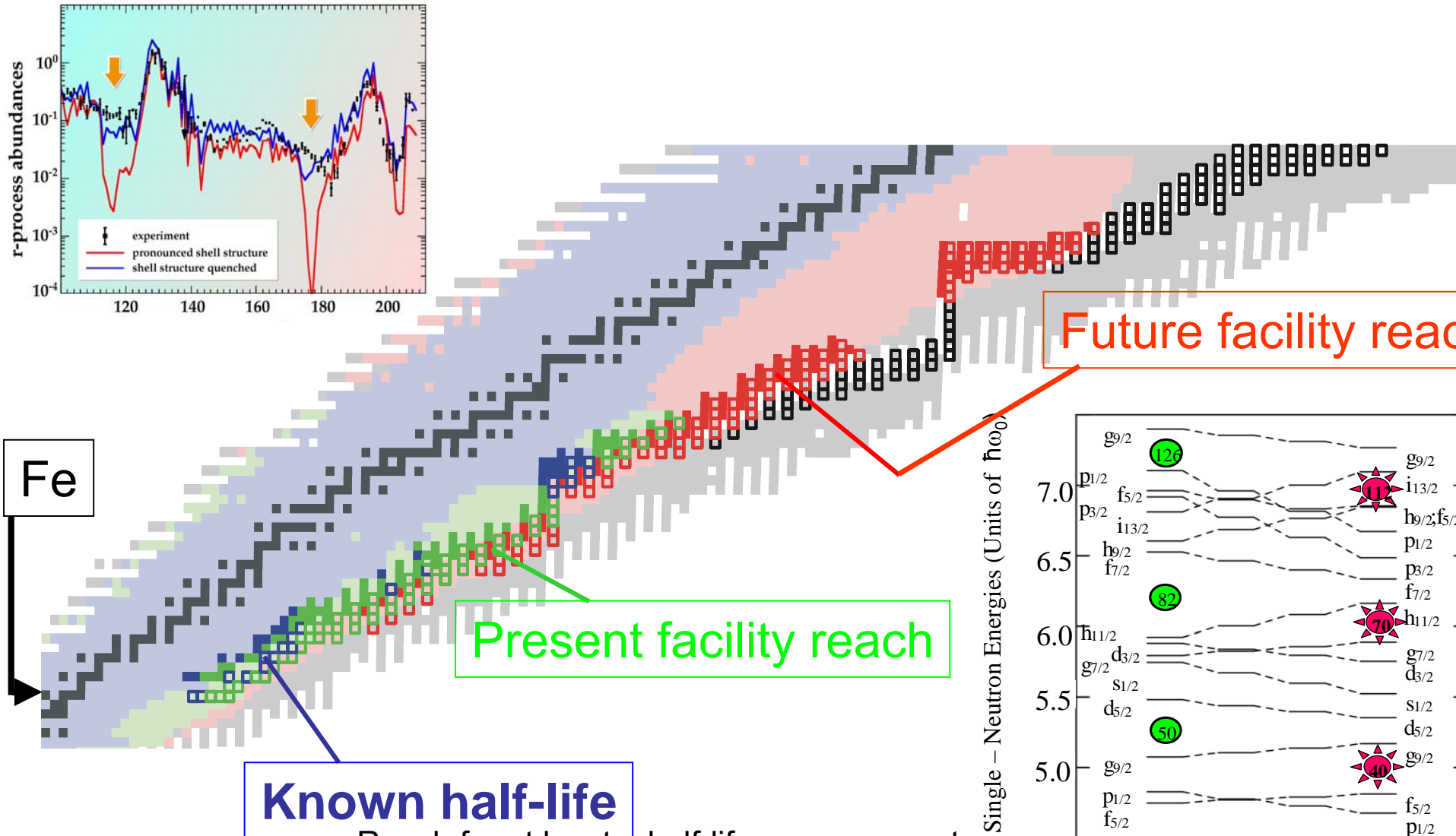


Figure 6: La/Eu ratios for a variety of metal-poor stars as a function of metallicity (24). The horizontal lines denote the total and r-process only solar system abundance ratios.

Science 299 (2003) 70



- Experimental data for extremely neutron-rich nuclei needed.
- Radioactive Ion Beam facilities
 - extension of the $N=50$ and 82 regions as close to drip-line as possible
 - $N=126$ “waiting-point” nuclei
- Ongoing sky surveys detect more metal-poor stars.
- More high-resolution optical spectroscopy in stars of varying metallicity is needed.
- More precise dating of the on-set of s-process nucleosynthesis.
- Scenarios/sites for the r-process
 - Effectively constrain “classical” parameter studies
 - Impact of the network calculations?
- Are there different r-processes?

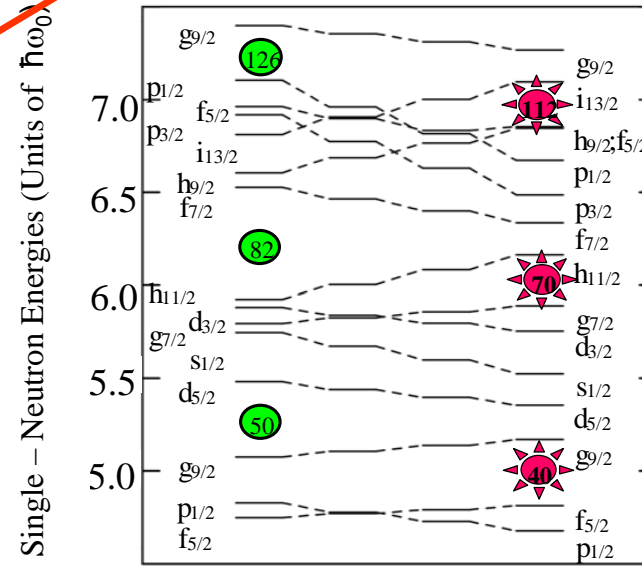


Future facility reach

Present facility reach

Known half-life

Reach for at least a half-life measurement



100% 70% 40% 10%

Strength of β -Term

new magic numbers

Fe